

3D relativistic hydro models for SS433: virtual views on precessing jets

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Outline

- 1 Introduction
- 2 Framework
- 3 Result - dynamics study
- 4 Result - radio study
- 5 Anexe

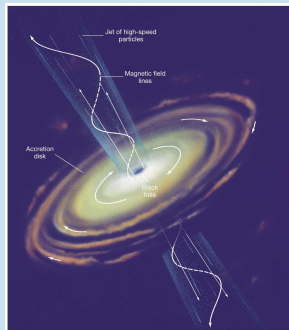
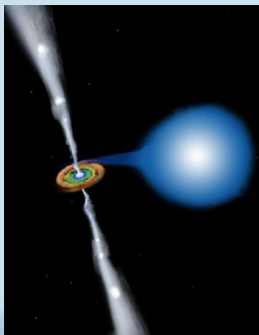
Team

- Monceau-Baroux Remi - Phd student at CmPA, KU Leuven
- Keppens Rony - Doctor at CmPA, KU Leuven - Supervisor
- Meliani Zakaria - Doctor at Luth, OBsPM - SRHD simulation of relativistic jet
- Porth Oliver - Doctor at CmPA, KU Leuven / Department of Applied Mathematics, The University of Leeds - Radio mapping

The presence of relativistic jets in the universe

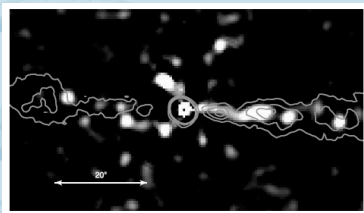
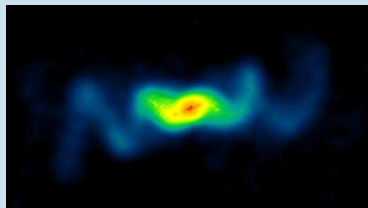
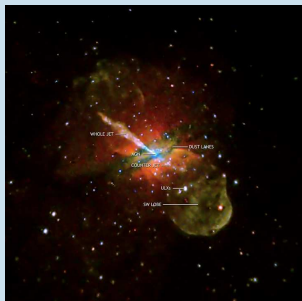
- AGN jets: FR-I, FR-II, etc: Expected to play an important role in the reheating of galaxy clusters. Ex: M87
The effect of angular opening on the dynamics of relativistic hydro jets - [Monceau-Baroux et al. 2012]
- X-ray binary jets. Ex: SS433
The effect of the Lorentz factor on the model of SS433 - [Monceau-Baroux et al. 2014 - DOI:10.1051/0004-6361/201322682]

False twins



- The accretion disk of a compact object (neutron star / black hole) releases matter in the form of winds. Many studies are made to explain how a jet can arise from these winds (see [Blandford & Payne 1982][Bogovalov and Tsinganos 2004])

Jets in the sky with diamants



Left top: Chandra X-ray Image of Centaurus A (Credit: NASA/CXC/CfA/R.Kraft et al.); Left bottom: 0.3 to 5.0 keV CHANDRA image of NGC 4261 (3C 270) after subtracting the diffuse component. The contours correspond to radio emission from a 4.9 GHz VLA observation (Zezas et al 2005); Right: VLA observations of SS433

To study or not study jets, why is it our question?

Aims for the study of X-ray binary associated jets:

- 1 Better understandings of relativistic jets:
 - 1 How does the precession of the jet affect the jet/medium interactions
 - 2 How do the properties of the jet (velocity/density ratio) affect the jet/medium interactions.
- 2 Comparison to observations, in case of SS433 with the VLA telescope. Need the ability to do radio mapping.

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How do we do that?

Framework

- 1 Relativistic version of hydrodynamical equations
- 2 Sygne gas equation of state
- 3 Code used for the simulations: MPI-AMRVAC
(<https://gitorious.org/amrvac/>)
- 4 Adaptive Mesh Refinement and Message Passing Interface

Where is Charly?

Problem?

What question should you ask yourself about that model?

Wait a minute!

Where is \vec{B} ?

Charly is not important...

Pro

Why \vec{B} shall be here: radio/X emissions = synchrotron

Con

Far away from the source: kinetic mechanisms should dominate magnetic ones. Actual measure of a few mG

For later on

Pressure taken as a proxy for magnetic field for radio emission.

What we know ...

Our input parameters are coming from observation:

- 1 The thermodynamic conditions of the ISM, pressure and density: P_{ISM}, ρ_{ISM} (Safi-Harb Oegelman 1997),
- 2 The energy flux of the jet, L_j (Brinkmann et al 2005),
- 3 The jet opening angle and the jet angle to its precession axis: α_j, θ_{prec} (Margon et al. 1979),
- 4 The velocity of the jet head: v_{head} (Roberts et al 2008).

... and what we do not!

We fix

- 1 The jet Lorentz factor, γ_j ,
- 2 $P_j = P_{ISM}$,

We need

- 1 The jet density, ρ_j .

Playing with known parameters

As for Meliani et al 2008 and Monceau et al 2012, we compute the integrated energy flux over the beam cross section as:

$$L_j = (\gamma_j h_j - 1) \rho_j \gamma_j \pi R_j^2 v_j, \quad (1)$$

where ρ_j , R_j , v_j are the jet density, radius and velocity.

$\rho_j h_j = \rho_j + \frac{\Gamma}{\Gamma-1} P_j$ is the enthalpy.

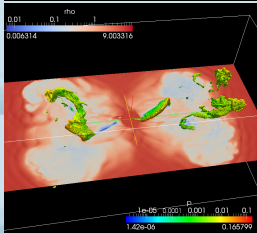
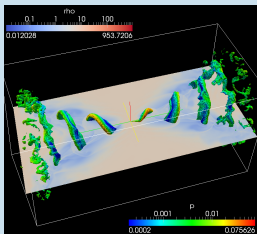
We can then obtain ρ_j .

4 cases for a global picture

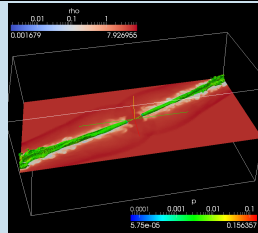
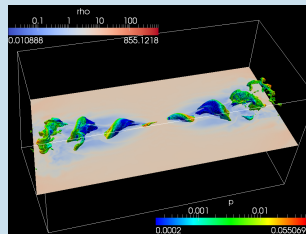
Case	γ_b (v_b)	η	θ_{prec}
A	1.036 (0.26c)	28.6	20°
B	1.87 (0.845c)	0.8	20°
C	1.036 (0.26c)	28.6	10°
D	1.87 (0.845c)	0.8	0°

Table : Parameters for the simulations. With $\eta = \gamma_j^2 \frac{\rho_j h_j}{\rho_{ISM} h_{ISM}}$ the inertia ratio.

Let's have a look



Top: case A, $\gamma = 1.036$ and $t = 2$,
 Bottom: case B, $\gamma = 1.87$ and $t = 2$,

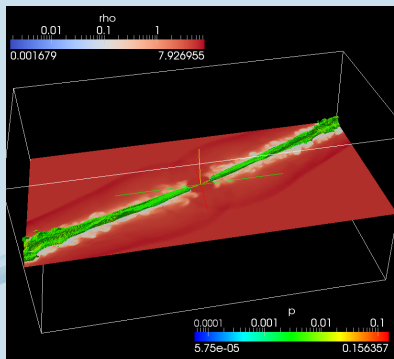


Top: case C, $\gamma = 1.036$ and $t = 2$,
 Bottom: case D, $\gamma = 1.87$ and $t = 0.5$.

Outline

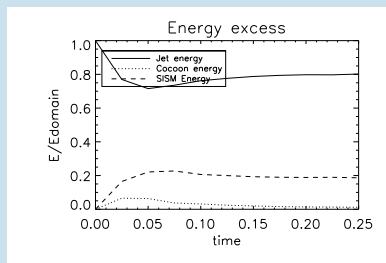
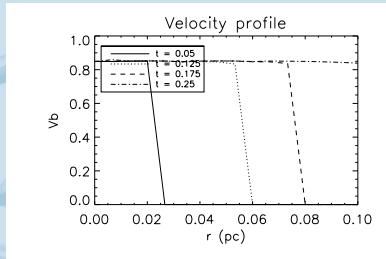
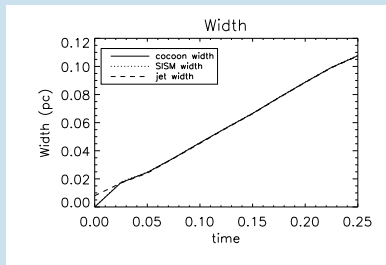
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Case D - Overview



- 1 'classic' static case
- 2 $\gamma = 1.87$, mildly relativistic
- 3 Bullet like propagation, canonical relativistic jet behavior

Case D - Dynamics

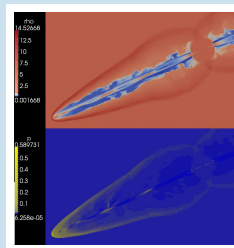
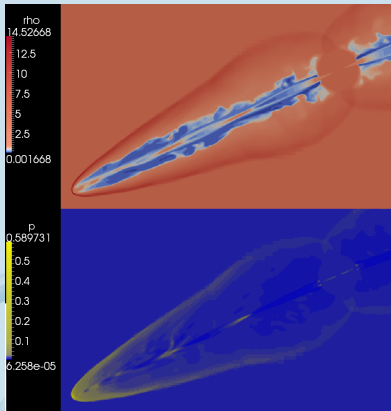


- 1 Path set by the jet head
- 2 It propagates with

$$v_j^{\text{mari}} = \frac{\sqrt{\eta}}{\sqrt{\eta+1}} v_{\text{beam}}$$
 with

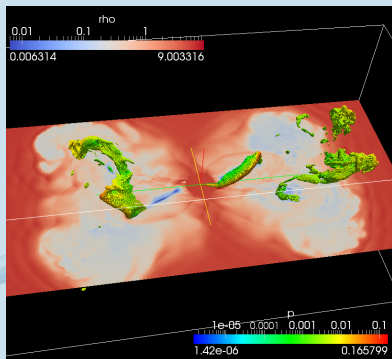
$$\eta = \gamma_j^2 \frac{\rho_j h_j}{\rho_{\text{ISM}} h_{\text{ISM}}}$$
- 3 'Near' flat velocity profile: only interaction at the head
- 4 Most energy in the jet beam: low interaction with the medium

Case D - internal structure



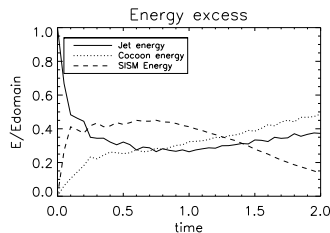
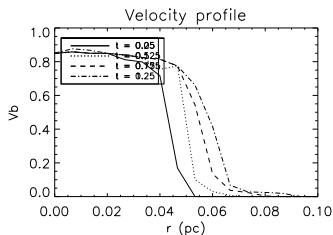
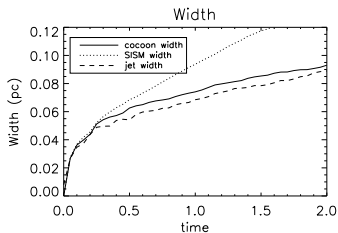
- ① Recollimation shocks
- ② Structured beam
- ③ Instabilities advected

Case B - Overview



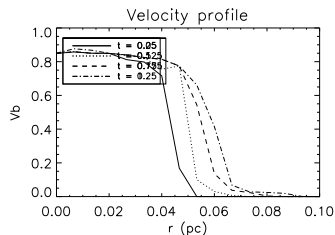
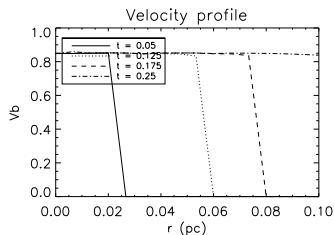
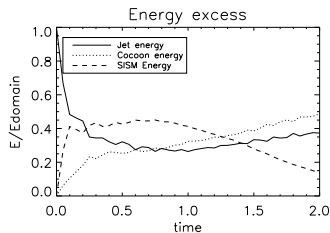
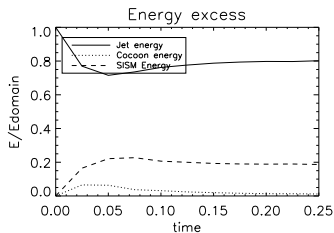
- 1 Precessing jet
- 2 Mildly relativistic
- 3 $\gamma = 1.87, \theta = 20^\circ$

Case B - Dynamics

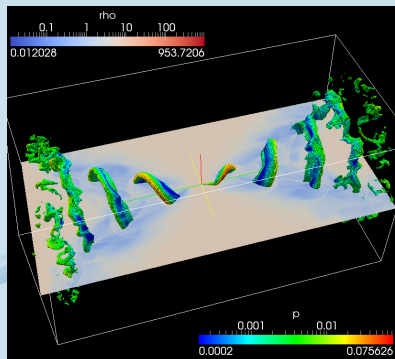


- 1 Deceleration of the jet head velocity to an asymptotic regime
- 2 Sub-sonic velocity of the jet head
- 3 Continuous deceleration along the path of the beam
- 4 Knee and ankle of the velocity profile
- 5 30% energy transferred to cocoon and 40% to the SISM

Effect of precession - case B and D

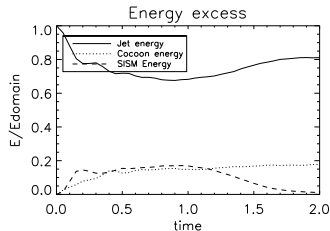
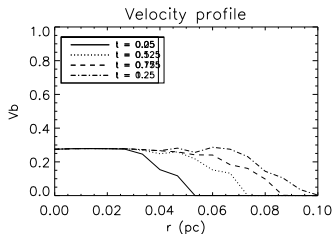
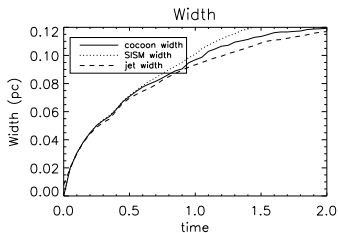


Case A - Overview



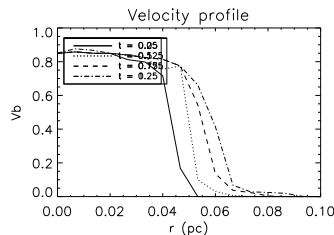
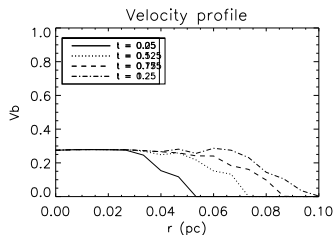
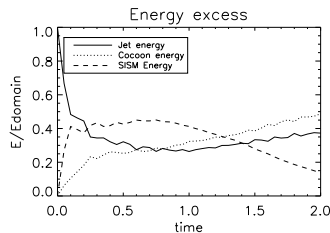
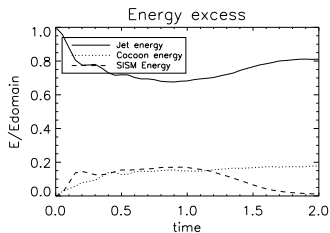
- 1 Canonical SS433 'kinematic model'
- 2 Barely relativistic
- 3 $\gamma = 1.036$, $\theta = 20^\circ$

Case A - Dynamics



- 1 Deceleration of the jet head velocity
- 2 Sub-sonic velocity of the jet head
- 3 Continuous deceleration along the path of the beam
- 4 40% energy transferred to cocoon and SISM (20% each)

Effect of Lorentz factor - case A and B



Conclusion on dynamics

- 1 Case D shows the formation of a structured beam and inner standing shocks known from the study of relativistic jets. It interacts weakly with the medium.
- 2 Precessing case where a shock propagates in front of the jet head display a knee and ankle velocity profile showing how the SISM is heated and accelerated by the shocks.
- 3 The precession increases the surface of interaction and the energy transfer
- 4 Increased Lorentz factor slows down the expansion of the interaction region. As the inertia ratio increases drastically with the Lorentz factor we observe the expected higher interaction with the medium.

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Radio Mapping

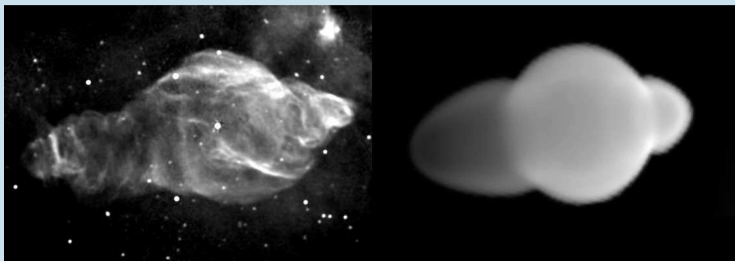
- Script for radio mapping by Oliver Porth (Porth 2014, DOI:10.1093/mnras/stt2176)
- We follow the evolution of the energy spectrum of the electrons: $f(\epsilon) = A\epsilon^{-\Gamma}$ for $\epsilon \leq \epsilon_\infty$, with $\Gamma = 0.6$.
- We use the emission equation from Camus et al (2009).

$$I = n_0 D^2 B_\perp \left(\frac{\rho_{e0}}{\rho_e}\right)^{-\frac{\Gamma+2}{3}} \epsilon^{1-\Gamma} \left(1 - \frac{\epsilon}{\epsilon_\infty}\right)^{\Gamma-2}, \quad (2)$$

where e and m are the electron charge and mass, ρ_e and ρ_{e0} are the electron density and initial density, ϵ_∞ the electron cut off energy, $\nu = c_1 B_\perp \epsilon^2$, $c_1 = 3e/4\pi m^3 c^5$, B_\perp is the component of the magnetic field normal to the line of sight in the fluid frame and $D = \nu_{obs}/\nu$ is the Doppler factor.

- B taken equal in intensity to $\sqrt{(P)}$
- Thin medium
- Ray tracing

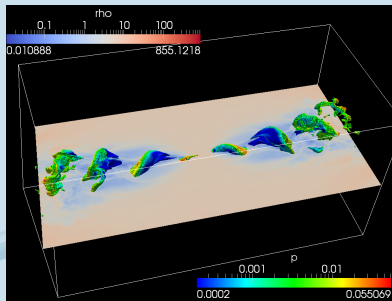
Zavalas simulation of SS433



Comparison of the radio-continuum image with the simulated column electronic density map of model M4. The left-hand panel shows the 1415-MHz image in grey-scale and contours of the W50 SNR, obtained with the VLA by Dubner et al. (1998), in equatorial coordinates (north is up). The right-hand panel shows the simulated map in a grey colour scale. A distance of 3 kpc to SS433 was assumed. (Zavala 2008)

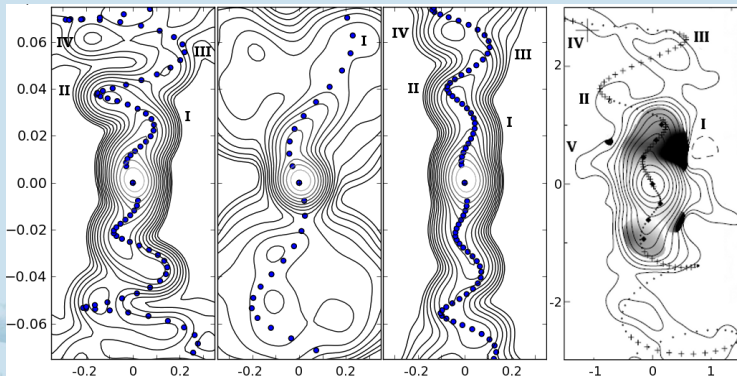
Need 10° precessing angle to reproduce the image at 20pc

Case C - Overview



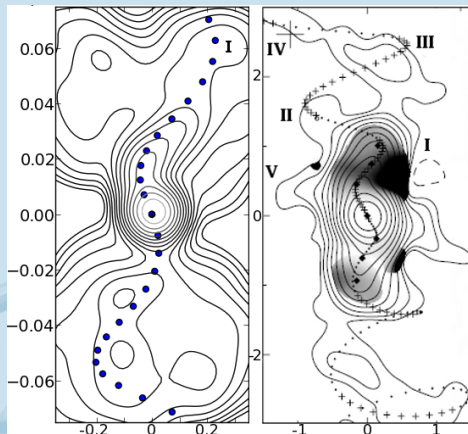
- 1 Barely relativistic
- 2 $\gamma = 1.036$, $\theta = 10^\circ$

Radio mapping - Overview



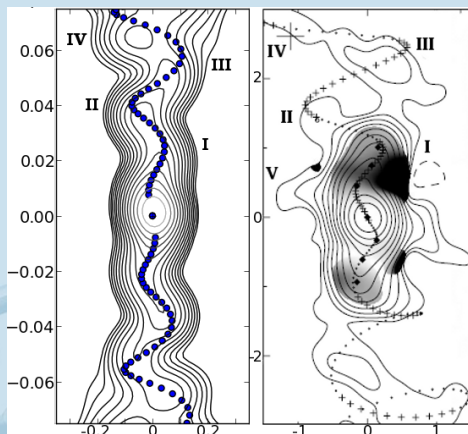
Left to right: Radio map from simulations Case A, Case B and case C. Units are in parsec, object is estimated to be at a distance of 5.5 kpc. All graphs overplot the kinematic model with parameters corresponding to the case. Right: VLA image of the microquasar SS433 in the constellation Aquila, adapted from Roberts et al. 2008, units are in arcsecond.

Case B - Too far



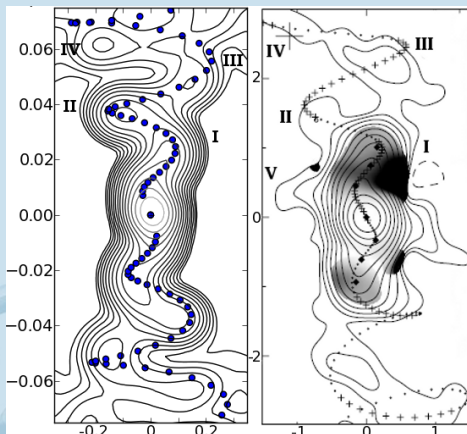
- 1 Radio elements too far from the source
- 2 Strong beaming effect

Case C - Too narrow



- 1 Radio elements too close to the precession axis.
- 2 Different precessing angle with time?

Case A - Good fit



- 1 Similar appearance
- 2 The kinematic model underestimates interactions on both simulations and observations.
- 3 Absence of the radio ruff

Conclusion on dynamics

- 1 Discrepancy at sub parsec scale and 20 parsec: time variation of the precessing angle? Recollimation?
- 2 Validation of the kinematic model for SS433. Only case A opening angle and Lorentz factor gets a similar picture to VLA observation.
- 3 The kinematic model needs to be corrected for interactions. It overestimates both simulations and observations.
- 4 Absence of the radio ruff: are they coming from the disk wind?

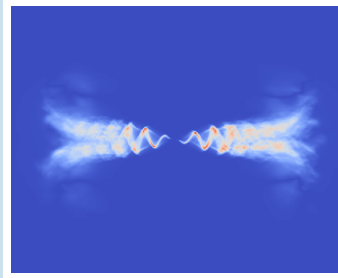
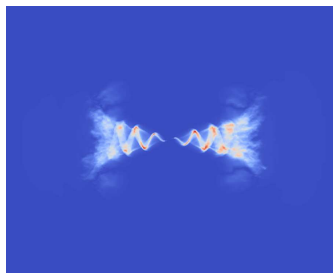
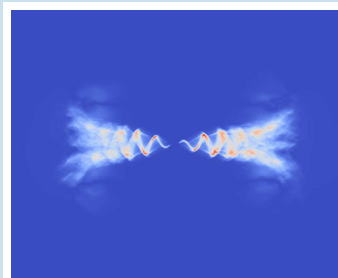
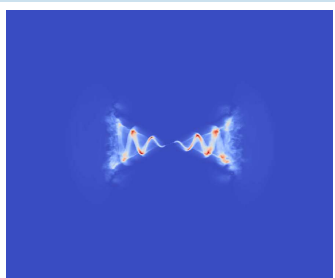
Wait a minute!

We have a problem:

- 1 Small scale (under parsec): 20°
- 2 Large scale (over 30 parsec): 10°



Spatial Evolution



Acknowledgement

Special thanks to Rony Keppens, Zakaria Meliani
and Oliver Porth
Dank u wel



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Equations of Relativistic Hydrodynamics

Equations of Relativistic Hydrodynamics

1 Continuity Equation

$$\frac{\partial \rho \gamma}{\partial t} + \vec{\nabla} \cdot \rho \gamma \vec{v} = 0 \quad (3)$$

2 Momentum Equation

$$\frac{\partial \vec{S}}{\partial t} + \vec{\nabla} \cdot (\vec{S} \vec{v}) + \vec{\nabla} p = 0 \quad (4)$$

3 Energy Equation

$$\frac{\partial \tau}{\partial t} + c^2 \vec{\nabla} \cdot (\vec{S} - \rho \gamma \vec{v}) = 0 \quad (5)$$

1 Momentum Density

$$\vec{S} = \frac{h}{c^2} \gamma^2 \rho \vec{v} \quad (6)$$

2 Specific Enthalpy

$$h = c^2 + \epsilon + \frac{p}{\rho} \quad (7)$$

3 Energy Density

$$\tau = \rho h \gamma^2 - p - \rho \gamma c^2 \quad (8)$$

Closure equation - The syngas equation of state

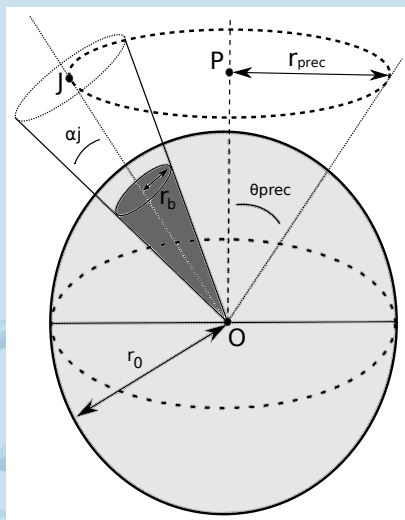
- 1 Mathews approximation to the Syngas equation

$$p = \left(\frac{\Gamma - 1}{2}\right)\rho\left(\frac{e}{m_p} - \frac{m_p}{e}\right) \quad (9)$$

- 2 Which gives a local effective polytropic index

$$\Gamma_{eff} = \Gamma - \frac{\Gamma - 1}{2}\left(1 - \frac{m_p^2}{e^2}\right) \quad (10)$$

Geometry



- 1 Binary system is not visible
- 2 Precession with 162 days
- 3 Overwrite central region
- 4 Double jet